TENSOR TRANSFORMATIONS (03)

I Main topics

- A Review of vector (first-order tensor) transformations
- **B** Second-order tensor transformations

Il Review of vector (first-order tensor) transformations

Consider a vector **T.** It has a magnitude and one associated direction (one subscript). If T transforms according to the "rules" of (3.1), then it is a first-order tensor.

Consider **T** in three coordinate systems: X, X', X".

$$T_{i}'' = a_{i''j'} T_{j'}$$
 and $T_{j}' = a_{j'k} T_{k}$ (3.1)

These can be expressed in matrix form too

$$\begin{bmatrix} T_{1"} \\ T_{2"} \\ T_{3"} \end{bmatrix} = \begin{bmatrix} a_{1"1} & a_{1"2} & a_{1"3} \\ a_{2"1} & a_{2"2} & a_{2"3} \\ a_{3"1} & a_{3"2} & a_{3"3} \end{bmatrix} \begin{bmatrix} T_{1'} \\ T_{2'} \\ T_{3'} \end{bmatrix} = \begin{bmatrix} a_{1'1} & a_{1'2} & a_{1'3} \\ a_{2'1} & a_{2'2} & a_{2'3} \\ a_{3'1} & a_{3'2} & a_{3'3} \end{bmatrix} \begin{bmatrix} T_{1} \\ T_{2} \\ a_{3'1} & a_{3'2} & a_{3'3} \end{bmatrix} \begin{bmatrix} T_{1} \\ T_{2} \\ T_{3} \end{bmatrix}$$

$$(3.2)$$

Substituting the expression for $T_{\dot{I}^{\,\prime}}$ on the right into the one of the left:

$$T_{i''} = a_{i''j'}a_{j'k}T_k$$
 (3.3)

$$\begin{bmatrix} T_{1"} \\ T_{2"} \\ T_{3"} \end{bmatrix} = \begin{bmatrix} a_{1"1'} & a_{1"2'} & a_{1"3'} \\ a_{2"1'} & a_{2"2'} & a_{2"3'} \\ a_{3"1'} & a_{3"2'} & a_{3"3'} \end{bmatrix} \begin{bmatrix} a_{1'1} & a_{1'2} & a_{1'3} \\ a_{2'1} & a_{2'2} & a_{2'3} \\ a_{3'1} & a_{3'2} & a_{3'3} \end{bmatrix} \begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix}$$

$$(3.4)$$

If the X" and X reference frames are the same, then T_i " = T_k

Let us replace each i" term in (3.3) by k:

$$T_k = a_{kj'} a_{j'k} T_k$$
 (3.5)

$$\begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} = \begin{bmatrix} a_{11'} & a_{12'} & a_{13'} \\ a_{21'} & a_{22}' & a_{23'} \\ a_{31'} & a_{32'} & a_{33'} \end{bmatrix} \begin{bmatrix} a_{1'1} & a_{1'2} & a_{1'3} \\ a_{2'1} & a_{2'2} & a_{2'3} \\ a_{3'1} & a_{3'2} & a_{3'3} \end{bmatrix} \begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} \tag{3.6}$$

This means that akj'aj'k must equal the identity matrix I:

$$\begin{bmatrix} a_{11'} & a_{12'} & a_{13'} \\ a_{21'} & a_{22'} & a_{23'} \\ a_{31'} & a_{32'} & a_{33'} \end{bmatrix} \begin{bmatrix} a_{1'1} & a_{1'2} & a_{1'3} \\ a_{2'1} & a_{2'2} & a_{2'3} \\ a_{3'1} & a_{3'2} & a_{3'3} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
(3.7)

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We can get this same result another way. We know that if we have three perpendicular unit vectors $\mathbf{n1}$, $\mathbf{n2}$, and $\mathbf{n3}$, with x, y, and z components, then

n1.n1 = 1, n1.n2 = 0, and n1.n3 = 0; n2.n1 = 0, n2.n2 = 1, and n2.n3 = 0;

n3.n1 = 0, n3.n2 = 0, and n3.n3 = 1.

These equations are expressed with matrices as:

So the sum of the squares of the direction cosines for a particular direction equals 1 (if the direction cosines for a given axis are dotted into themselves, then the answer is one; the directions are parallel). If the direction cosines for perpendicular directions are dotted into each other, the result is zero.

Replacing x with 1', y with 2', and z with 3'

$$\begin{bmatrix} a_{11'} & a_{12'} & a_{13'} \\ a_{21'} & a_{22'} & a_{23'} \\ a_{31'} & a_{32'} & a_{33'} \end{bmatrix} \begin{bmatrix} a_{11'} & a_{21'} & a_{31'} \\ a_{12'} & a_{22'} & a_{32'} \\ a_{13'} & a_{23'} & a_{33'} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
(3.9)

Because $\cos \theta_{ij} = \cos \theta_{ji}$, we can switch the order of the subscripts on the terms in the second matrix:

$$\begin{bmatrix} a_{11'} & a_{12'} & a_{13'} \\ a_{21'} & a_{22'} & a_{23'} \\ a_{31'} & a_{32'} & a_{33'} \end{bmatrix} \begin{bmatrix} a_{1'1} & a_{1'2} & a_{1'3} \\ a_{2'1} & a_{2'2} & a_{2'3} \\ a_{3'1} & a_{3'2} & a_{3'3} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
(3.10)

This matches (3.7).

III Second-order tensor transformations

Consider two vectors u_i and v_k . Then if their components are transformed into different reference frames:

$$u_{i}' = a_{i'i}u_{i}$$
 and $v_{\ell}' = a_{\ell'k}v_{k}$ (3.11)

Suppose we multiply each component of $u_{\bar{l}}$ by each component of v_{ℓ} to give a 3x3 matrix:

$$u_{j'}v_{\ell'} = (a_{j'i}u_i)(a_{\ell'k}v_k) = a_{j'i}a_{\ell'k}u_iv_k$$
(3.12)

Now let the 3x3 matrix product uv = A

$$A_{j'\ell'} = a_{j'i}a_{\ell'k}A_{ik}$$
 (3.13)

$$[A_{j'\ell'}] = [a_{j'i}][A_{ik}][a_{\ell'k}]$$
 (3.14)

Quantities which transform this way are called second-order tensors. The components have magnitudes and two associated directions (2 subscripts). Stress components transform like this, so stress is considered to be a second-order tensor.