

## SUBSIDENCE MECHANICS: HEAT FLOW ANALOG (38)

### I Main Topics

- A Motivation: Why investigate heat flow?
- B Development of 1-D heat flow equation as analog for consolidation
- C Finite-difference interpretation of heat flow equation
- D Dimensional analysis

### II Motivation

- A Heat flow equation has the same form as the consolidation equation but is easier to grasp
- B Diffusion of heat analogous to diffusion of excess pore pressure
- C Many analytic solutions for heat flow (e.g., Carslaw and Jaeger, 1984)
- D Many analogous equations of great use

## E Flow analogs

Flowing quantity	Incompressible Fluid	Heat	Chemical Species
Conserved quantity	Mass	Heat Energy	Molecules
1-D flux law	Darcy's law $q = -k \frac{\partial H}{\partial x}$	Fourier's law $q = -k \frac{\partial T}{\partial x}$	Fick's law $J = -D \frac{\partial c}{\partial x}$
Flux term	$q = \text{volume flux density}$ $\text{m}^3/(\text{m}^2 \cdot \text{sec})$	$q = \text{heat flux density}$ $\text{joules}/(\text{m}^2 \cdot \text{sec})$	$J = \text{diffusion flux}$ $\text{moles}/(\text{m}^2 \cdot \text{sec})$
Coefficient	$k = \text{hydraulic conductivity}$ $\text{m/sec}$	$k = \text{thermal conductivity}$ $\text{Joules}/(\text{m} \cdot \text{K} \cdot \text{sec})$	$D = \text{diffusivity}$ $\text{m}^2/\text{sec}$
Potential term	$H = \text{head (m)}$	$T = \text{temperature (\text{K})}$	$c = \text{concentration (moles/m}^3)$
1-D diffusion law	$\frac{\partial H}{\partial t} = \alpha \frac{\partial^2 H}{\partial x^2}$	$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$	$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2}$
Coefficient	$\alpha = \text{hydraulic diffusivity}$ $\text{m}^2/\text{sec}$	$\alpha = \text{thermal diffusivity}$ $\text{m}^2/\text{sec}$	$D = \text{diffusivity}$ $\text{m}^2/\text{sec}$
Steady state flow (Term on left side of diffusion law = 0)	$\nabla^2 H = 0$	$\nabla^2 T = 0$	$\nabla^2 c = 0$

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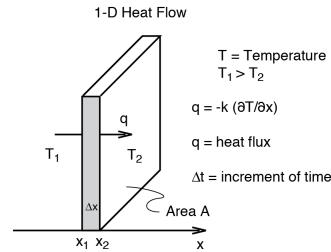
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## III Development of 1-D heat flow equation as analog for consolidation

### A Isotropic, uniform material

### B Definition of terms

- 1  $U = \text{heat energy (joules)}$
- 2  $x = \text{position (meters)}$
- 2  $t = \text{time (seconds)}$
- 3  $q = \text{heat flux (joules/(meter}^2 \text{ sec)})$ 
  - a Rate of heat energy transfer per unit area per unit time
  - b Heat flux can vary with time and position, so  $q = q(x,t)$
- 4  $T = \text{temperature (\text{)} }$   
Temperature can vary with position and time, so  $T = T(x,t)$



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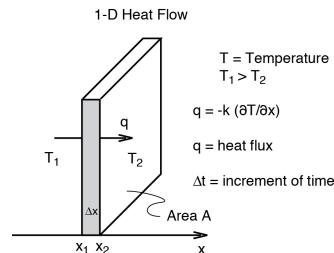
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## C Fourier's Law of Heat Conduction (1-D)

$$q = -k \frac{\partial T}{\partial x}$$

- 1 Heat flow ( $q$ ) scales with the temperature gradient ( $\partial T / \partial x$ )
- 2  $k$  = coefficient of thermal conductivity
  - a Dimensions: Joules sec<sup>-1</sup> m<sup>-1</sup> K<sup>-1</sup>
  - b  $k$  assumed to be constant
- 3 Dimension check

$$\frac{\text{Joules}}{\text{m}^2 \text{ sec}} = \frac{\text{Joules}}{\text{m sec}^\circ \text{K}} \frac{^\circ \text{K}}{\text{m}}$$



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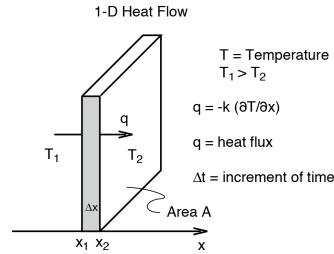
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## C Fourier's Law of Heat Conduction (1-D)

$$q = -k \frac{\partial T}{\partial x}$$

- 4 The minus sign
  - a For heat to flow from x<sub>1</sub> to x<sub>2</sub>, where x<sub>1</sub> < x<sub>2</sub>, T(x<sub>1</sub>) > T(x<sub>2</sub>).
  - b Positive heat flow corresponds to a drop in temperature, requiring k to be negative
- 5 Partial derivative used because T is a function of x and t.
- 6 Finite difference approximation:

$$q = -k \frac{\Delta T}{\Delta x}$$



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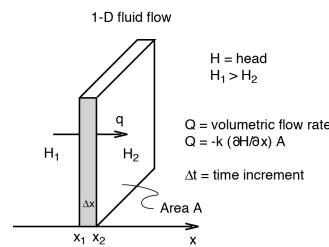
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## D Fluid flow analog (slow laminar flow)

$$q = -k \frac{\partial H}{\partial x}$$

- 1 Volumetric flux ( $q$ ) scales with the head gradient ( $\partial H / \partial x$ )
- 2  $k$  = hydraulic conductivity
  - a Dimensions: m/sec
  - b  $k$  assumed to be constant
  - c  $k$  depends on the intrinsic permeability of the material, the degree of saturation, and on the density and viscosity of the fluid
- 3 Dimension check

$$\frac{m^3}{m^2 \text{ sec}} = \frac{m}{\text{sec}} \frac{m}{m}$$



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## E Heat Flow Equation (Conservation of energy)

- Change in heat energy = heat in – heat out

$$\Delta U_{heat} = (\Delta T)(mass)(specific\ heat)$$

$$\Delta U_{heat} = (A)(\Delta t)[q(x=x_1) - q(x=x_2)]$$

$$(\Delta T)(mass)(specific\ heat) = (A)(\Delta t)[-Δq]$$

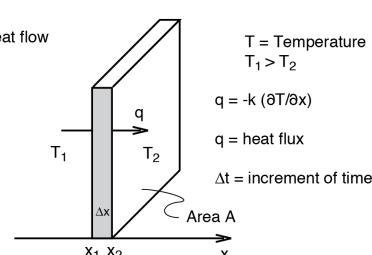
$$(\Delta T)(ρAΔx)(c) = (A)(\Delta t)\left[-\frac{\partial q}{\partial x}\Delta x\right]$$

$$\frac{(\Delta T)(ρ)(c)}{(\Delta t)} = \left[ -\frac{\partial}{\partial x} \left( -k \frac{\partial T}{\partial x} \right) \right]$$

q stated using Fourier's Law

$$\frac{(\Delta T)(ρ)(c)}{(\Delta t)k} = \frac{\partial^2 T}{\partial x^2}$$

Heat Flow



$$\frac{\partial T}{\partial t} = α \frac{\partial^2 T}{\partial x^2}$$

α = thermal diffusivity  
= k/ρc

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## F Heat Flow Equation

- 1-D form:  $K \frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2}$  parabolic differential equation
- 2-D form:  $K \frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2}$
- 3-D form:  $K \frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2}$
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## G Laplace equation

- Applies to steady state distribution of temperature
- Temperature does not change as a function of time
- $\frac{\partial T}{\partial t} = 0$
- 1-D:  $\frac{\partial^2 T}{\partial x^2} = 0$
- 2-D:  $\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0$
- 3-D:  $\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} = 0$
- General:  $\nabla^2 T = 0$

$$\text{Curvature}(T_{1-D}) = \frac{\frac{d^2 T}{dx^2}}{\sqrt{1 + \left(\frac{dT}{dx}\right)^2}}^{3/2}, \text{ so } \frac{d^2 T}{dx^2} = 0 \text{ means curvature}(T_{1-D}) = 0$$

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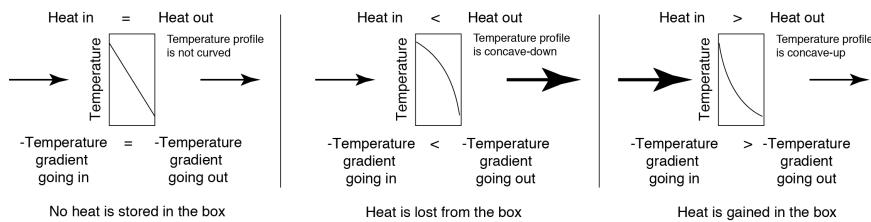
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## H Relationship between 1-D temperature profile and heat change

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

- To a good approximation, the rate of temperature change with time scales with the curvature of the temperature profile
- If the 1-D temperature profile isn't curved, then no change in heat energy occurs in the slab (i.e., steady state exists)



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## I Fluid Flow Equation

(Conservation of mass for incompressible fluid)

- Change in fluid mass = mass in – mass out

$$\Delta \text{mass} = \left( \frac{\Delta \text{volume}}{\Delta H} \right) (\text{density})(\Delta H)$$

$$\Delta \text{mass} = (A)(\Delta t)(\text{density})[q(x=x_1) - q(x=x_2)]$$

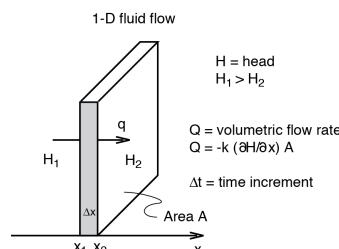
$$\left( \frac{\Delta \text{volume}}{\Delta H} \right) (\text{density})(\Delta H) = (A)(\Delta t)(\text{density})[-\Delta q]$$

$$\left( \frac{\Delta \text{volume}}{A \Delta H} \right) (\Delta H) = (\Delta t) \left[ \frac{-\partial q}{\partial x} \Delta x \right]$$

$$S \frac{(\Delta H)}{(\Delta t)} = \left[ -\partial \left( -k \frac{\partial H}{\partial x} \right) \right]$$

S = storativity

$$S \frac{\Delta H}{\Delta t} = k \frac{\partial^2 H}{\partial x^2}$$



$$\frac{\partial H}{\partial t} = \alpha \frac{\partial^2 H}{\partial x^2}$$

$\alpha$  = hydraulic diffusivity  
 $= k/S$

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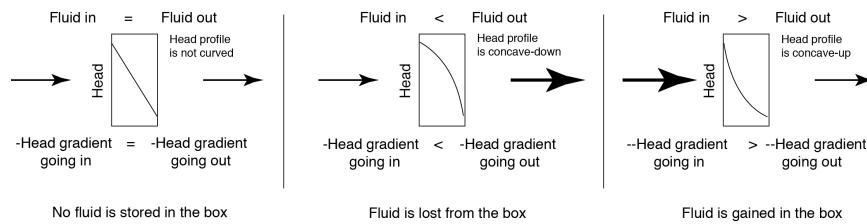
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## J Relationship between head profile and fluid volume change

$$\frac{\partial H}{\partial t} = \alpha \frac{\partial^2 H}{\partial x^2}$$

- To a good approximation, the rate of fluid content change with time scales with the curvature of the head profile
- If the head profile isn't curved, then no change in fluid content occurs in the slab



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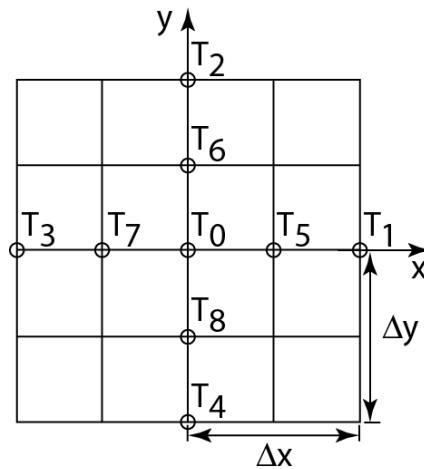
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## IV Finite difference interpretation of heat flow equation

- $\nabla^2 T = 0$
- The value of  $T$  (here  $T$  = temperature) at a given point is the average of the values at the nearest neighboring points on a square grid (see notes on wave eqn)

$$T_0 = \frac{T_1 + T_2 + T_3 + T_4}{4}$$



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## V Dimensional analysis

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

- A Question: How does the thickness  $H$  of a plate control the time the plate takes to cool?
- B Consider the dimensions of the terms in the heat flow equation and the plate thickness  $H$

$[t_c] = \text{cooling time}$

$[T] = {}^\circ K$

$[\alpha] = (\text{length})^2 (\text{time})^{-1}$

$[H] = [x] = \text{length}$

## V Dimensional analysis

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

- C Now consider the cooling time at a point in the plate
- D The equation for the cooling time  $t_c$  must be dimensionally consistent, and can only depend on the relevant factors. So

$$t_c = CT^a \alpha^b H^c, \quad [t_c] = \text{time} \quad [T] = {}^\circ K \\ [\alpha] = (\text{length})^2 (\text{time})^{-1} \quad [H] = \text{length}$$

where  $C$  is an unknown dimensionless constant

Hence  $[t_c] = [T]^a [\alpha]^b [H]^c$

$$(\text{time})^1 = ({}^\circ K)^a \left( \frac{\text{length}^2}{\text{time}} \right)^b (\text{length})^c$$

$$(\text{time})^1 = ({}^\circ K)^a (\text{length})^{2b+c} (\text{time})^{-b}$$

E By inspection

$$a = 0; b = -1; c = 2$$

So

$$t_c = C \alpha^{-1} H^2$$

## V Dimensional analysis

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

F Meaning of solution  $t_c = C\alpha^{-1}H^2$

- 1 The time for cooling the plate to some fraction of its initial temperature is proportional to  $1/H^2$
- 2 Doubling the plate thickness quadruples the cooling time
- 3 If  $C \approx 1$ , then  $t_c \approx \frac{H^2}{\alpha}$

This can be used for rough estimates of the cooling time

## References

- Carslaw, H.S., and Jaeger,J.C., 1984,  
Conduction of heat in solids: Clarendon Press,  
Oxford, 510 p.
- Wang, H.F., and Anderson, M.P., 1982,  
Introduction to groundwater modeling:  
Academic Press, San Diego, 237 p.