The average vertical mixing coefficient for the oceanic thermocline

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ABSTRACT

The new estimate of the average vertical mixing coefficient for the oceanic thermocline, based on all GEOSECS tritium data, is $1.7 \pm 0.2 \, \text{cm}^2 \, \text{s}^{-1}$.

1. Introduction

Until a reliable general circulation model of the ocean has been developed, those interested in modeling the uptake of fossil fuel CO$_2$ by the sea and heat buffering exerted by the sea will have to depend on simple transfer models calibrated using radioisotope data. To date, we have only three isotope data sets suitable for global modeling, the distributions of natural radiocarbon and of bomb produced radiocarbon and tritium (at the time of the GEOSECS surveys; Atlantic 1972–73; Pacific 1973–74; and Indian 1978). However, the natural pre-bomb radiocarbon data are limited and the GEOSECS radiocarbon data need to be corrected for the pre-bomb radiocarbon contribution. Therefore, the tritium distribution holds the most valuable information. Also the time scale for the penetration of tritium is nearly matched to the time scales of the CO$_2$ and heat transients. Thus it is important that the results of the GEOSECS surveys of tritium be summarized in a way that make them useful to those conducting such modeling efforts.

Our approach was to obtain area-averaged depth distributions for tritium in various regions of the ocean. We then sought that mathematical function which would best fit these distributions. It turned out that a Fickian diffusion law is appropriate for this purpose. Because of this, the most convenient means of expressing the tritium distribution is in terms of vertical mixing coefficients with the same unit of normal diffusion coefficient ($\text{cm}^2 \, \text{s}^{-1}$). While these vertical mixing coefficients have no direct physical meaning in that tritium is thought to move into the body of the sea along isopycnal horizons, they do permit the mean distribution of tritium in the sea as a whole or in a given region of the sea to be reproduced. We also realize that these mixing coefficients are probably applicable only to penetration times close to that for the mean age of ocean tritium as of the GEOSECS surveys. Further, we realize that because of differing surface boundary conditions, these mixing coefficients may not apply equally well to CO$_2$ and heat even for the same penetration time scales. However until more transient tracer data is in hand allowing us to extend the penetration time scale and to carefully intercompare the distributions of various tracers (i.e. tritium, $^{85}$Kr and freons), this is the best approximation we can make.

In the one-dimensional box-diffusion model of Oeschger et al. (1975), the well-mixed surface ocean layer is coupled to a deep ocean layer through which material (e.g. fossil CO$_2$, bomb $^{14}$C, $^3$H, etc.) is transported by finite vertical mixing processes. By fitting the average vertical profile of pre-bomb $^{14}$C in the ocean to the model, Oeschger et al. (1975) estimated that the average vertical mixing coefficient ($K$) in the deep layer is most
likely about 1.26 cm$^2$ s$^{-1}$ with an acceptable range of from 0.95 to 1.78 cm$^2$ s$^{-1}$. As mentioned earlier and cautioned by Broecker et al. (1980), the vertical mixing coefficient $K$ is a purely empirical parameter which takes into account all the processes that transfer tracers across density horizons in the oceanic thermocline, e.g. vertical eddy mixing, and advection along isopycnal-horizons which outcrop at surface ocean.

Using the box-diffusion model of Oeschger et al. (1975) and the tritium concentrations of rain waters at 50°N as the time-dependent tritium input function (Weiss and Roether, 1980), Broecker et al. (1980) first established the relationship between the mean penetration depth of tritium ($= \int_0^z C \, dz/C_0$, where $C_0$ and $C$ are the concentration of tritium at surface and depth $z$ respectively) and $K$. Then, from the estimation of the global mean penetration depth of tritium from GEOSECS Atlantic and Pacific data, they obtained a $K$ of 1.7 cm$^2$ s$^{-1}$ which is within the range of Oeschger et al. (1975). Broecker et al. (1980) also obtained a $K$ in the range of from 1.9 to 3.3 cm$^2$ s$^{-1}$ using the GEOSECS Atlantic and Pacific 14C data. However, considering the large uncertainty involved in subtracting the pre-bomb 14C from the observed 14C data, the $K$'s estimated from bomb 14C have a large uncertainty.

Based on the GEOSECS tritium data from the northern hemisphere, Vicceli et al. (1981) proposed a $K$ value of 5.1 cm$^2$ s$^{-1}$, which is four times greater than that obtained by Oeschger et al. (1975). Vicceli et al. (1981) treated the bottom tritium maxima in the North Atlantic GEOSECS stations as cross isopycnal vertical mixing inputs from the warm surface in their model fit (instead of large-scale lateral advective inputs from the source areas of the North Atlantic Deep Water). This may explain, in part, their large $K$ value.

Considering the importance of the $\tilde{K}$ value in predicting the future atmospheric CO$_2$ concentration (Oeschger et al. 1975) and the global sea level changes (Gornitz et al., 1982), we present here a new estimate of $\tilde{K}$, based on all GEOSECS tritium data (Östlund et al., 1976, 1979, 1980) and various approaches. For even geographical coverage in the Atlantic Ocean, the tritium data from stations 3B, 8, 28C and 31B of German cruise GS7309 and 30, 45, 51 of German cruise GS7205 (W. Roether, personal communication) were also included.

2. Results and discussion

The ocean is divided into nine zones as shown in Fig. 1. The average concentration profiles of tritium in each zone (Fig. 1) were obtained by the following procedures: the tritium profiles in each 10° or 20° latitudinal strip were averaged first, then the averages of latitudinal strips within each zone were averaged again but using the area of each latitudinal strip as a weighting factor. In averaging the tritium profiles from the north Atlantic zone (20°N to 70°N), the advective inputs of tritium at the bottom of GEOSECS stations 3, 5, 11 and 23 were subtracted by assuming an exponential decrease of tritium concentrations at depths below 1500 m.

The solid line corresponding to a given surface mixed layer thickness ($h$) and vertical mixing coefficient ($K$) in Fig. 1 represents a good model fit to the averaged tritium profile in each zone, using the box-diffusion model of Oeschger et al. (1975). If the mixed layer thickness $h$ is increased (or decreased) by 25 m, one still can obtain a reasonably good fit in Fig. 1 by decreasing (or increasing) the K value by about 0.1 to 0.2 cm$^2$ s$^{-1}$, except the profile for the North Atlantic. The tritium concentrations of rain water at 50°N and 50°S were used as the tritium input function for northern zones (>20°N) and southern zones (>20°S), respectively (Weiss and Roether 1980). Both time-dependent tritium input functions were expressed as fractions of the total rain inputs between 1952 and 1977 for each hemisphere. For the equatorial zones (20°N to 20°S), we used the average of 50°N and 50°S. The model ran from 1952 to 1973 for the Atlantic, to 1974 for the Pacific and to 1977 for the Indian Ocean. The surface tritium concentration of the model-produced tritium profiles was always normalized to the “observed” tritium concentration of the surface mixed layer. Therefore, we needed to know only the relative magnitude of the time-dependent tritium input function.

The good fit of the model curves to the “observed” data points (Fig. 1) indicates that the tritium transport in the surface 2000 m of the ocean can be represented by a Fickian-type diffusion process on a regional scale.

The global average of the vertical mixing coefficient of the ocean ($\tilde{K}$) can be estimated from the relationship:

$$\tilde{K} = (\sum \sqrt{K_if_i})^2,$$  

(1)
Fig. 1. The box-diffusion model fits of the average tritium concentration profiles in various regions.

where $K_i$ and $f_i$ are the vertical mixing coefficient and the fraction of the total ocean area in zone $i$, respectively (Viecelli et al., 1981). The area-weighted average of the square root of the regional vertical mixing coefficients as the global average is justified, since the integrated uptake of tritium for any region is proportional to the square root of the vertical mixing coefficient (Viecelli et al., 1981). Using the $K_i$ values in Fig. 1, we obtain the $\bar{K}$ value of $1.7 \pm 0.2$ cm$^2$ s$^{-1}$ (with the average surface
mixed-layer thickness of about 75 m) which is consistent with the estimates by Oeschger et al. (1975) and Broecker et al. (1980). The uncertainty of the $K$ value is mainly contributed by the uncertainty in the surface mixed-layer thickness of $\pm 25$ m. The area-weighted average of the surface mixed-layer thickness over the whole ocean is estimated independently to be about 70 m using monthly temperature and salinity profile data of several hundred stations (Inez Fung, personal communication).

Another way to estimate $\bar{K}$ is to normalize the tritium profiles to the surface concentration in each zone, average the normalized tritium profiles using the area of each zone as a weighing factor, and fit the global average to the box-diffusion model using the average of the tritium input functions of rain water at 50°N and 50°S as in the cases of the equatorial regions (Fig. 2). The best fit $\bar{K}$ is again $1.7 \pm 0.2 \text{ cm}^2 \text{s}^{-1}$ with $h = 75$ m, if the model is run from 1952 to 1974. If the model is run from 1952 to 1973 or to 1975, $\bar{K}$ is $1.8 \pm 0.2$ or $1.6 \pm 0.2 \text{ cm}^2 \text{s}^{-1}$.

As of 1973, the total tritium advected into the bottom water of the north Atlantic (20°N to 70°N) was estimated to be about 13% of the surface tritium (0–2000 m) in the same zone. If we assume that an equal amount of tritium was also advected into the Arctic bottom water, then it is equivalent to an increase in the effective surface area of the north Atlantic zone of 26%. Therefore, by eq. (1) the effective $\bar{K}$ can be about $1.8 \pm 0.2 \text{ cm}^2 \text{s}^{-1}$.

We have also estimated $\bar{K}$ by the average mean penetration depths of tritium from three oceans. We first average the mean penetration depths of tritium for GEOSCECS stations within a 10° or 20° latitudinal strip in each ocean. Then the averages of latitudinal strips in each oceanic hemisphere (north and south) were averaged again using the area of each latitudinal strip as a weighting factor. The results are summarized in Table 1. The tritium concentrations usually decrease rapidly with depth and oscillate around zero with the standard error of about $\pm 0.06$ to $\pm 0.09$ T.U. below a certain depth when a 0.02 T.U. blank correction (Östlund et al., 1980) is applied to all GEOSCECS tritium data. Therefore, the depth interval(s) where the tritium concentrations oscillate around the standard error (indicating near-zero concentration) are not included in the mean penetration depth calculation for each GEOSCECS station. The relationships between $K$ and the mean penetration depth for each oceanic hemisphere (Fig. 3) were obtained by using the box-diffusion model with the tritium concentration of rain water at 50°N and 50°S as
the tritium input functions for northern and southern hemispheres, respectively. The average surface mixed-layer depth was taken to be 75 m (Fig. 2). As before, the model ran from 1952 to 1973 for the Atlantic, to 1974 for the Pacific and to 1977 for the Indian Oceans. The $K$'s for corresponding mean penetration depth in each oceanic hemisphere, as obtained from Fig. 3, are given in Table 1. The $\bar{K}$ is $1.8 \pm 0.2 \text{ cm}^2\text{s}^{-1}$ if $K$'s from Table 1 are applied to eq. (1). If the bottom tritium inputs in the North Atlantic GEOSECS stations are subtracted from the mean penetration depth calculation, $\bar{K}$ is reduced to $1.7 \pm 0.2 \text{ cm}^2\text{s}^{-1}$.

In short, our refined estimate of $\bar{K}$ is within the range of previous estimates by Oeschger et al. (1975) and Broecker et al. (1980). Therefore, their calculations with regard to the fossil CO$_2$ uptake by the oceanic thermocline are still appropriate.

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LETTER TO THE EDITOR

A reply to the paper by Austin and Green

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We would like to clarify that the mixed-layer depth ($h$) and the vertical coefficient ($K$) in our paper (Li et al. (1984)) were obtained simultaneously by curve fitting. Our purpose was not to elucidate actual mixing mechanisms taking place in the ocean but rather to provide a practical interim means of calculating CO$_2$ uptake, heat buffering, pollutant dispersal ... in the sea. We observed that the shapes of horizontally averaged tritium profiles for large regions of the sea could be adequately approximated by the diffusion of tritium from a well-mixed reservoir into a semi-infinite half space. We were very careful to state that we do not believe that vertical eddy mixing was responsible for the actual oceanic transport. Thus the Austin and Green criticism is misdirected. Contrary to their implication, we used no oceanographic information to obtain our value of $h$ (the mixed-layer depth). It was chosen to fit the shape of the horizontally averaged tritium profiles.

We have emphasized that our $K$ is a purely empirical parameter to reproduce the mean distribution of tritium in the surface ocean (0-2000 m). Our $K$ takes into account all the physical processes which transfer tritium across density horizons in the oceanic thermocline. Therefore, the mixing due to baroclinic eddies is included in addition to mixing and advection along isopycnal horizons which outcrop at surface ocean. The ventilation of thermocline through outcrops of isopycnal horizons is well documented by tritium and $^{3}$He data (Jenkins, 1980). However, it is only one of many possible tritium transport mechanisms and has nothing to do with our model calculation of $K$ and $h$ as suggested by Austin and Green. We have no argument with the obvious temporal and spacial variations of $h$ and $K$ in the ocean suggested by Austin and Green.

The very reason why we averaged the tritium profile data in a large regional scale is to smooth out the subsurface tritium maxima which often occur in individual profiles and which cannot be fitted by a simple one-dimensional representation. In order to fit the tritium profile data in a finer scale, one needs at least a two-dimensional diffusion-advection model. It was not our intention to make such a model.

We admit that our model fit to the north Atlantic data is less than perfect, especially with regard to the slight tritium minimum around 700-900 m. The tritium minimum was caused by the northward penetration of the high silica but low tritium antarctic circumpolar intermediate water (Broecker and Peng, 1982, Figs. 1-14). However, that our $K$ for the north Atlantic is actually a maximum estimation is clear from our Fig. 1. We define the mixed-layer depth as the depth of the well-mixed surface water with near constant salinity and tritium concentration. Since the tritium concentration in the North Atlantic decreases with depth beneath the base of our mixed layer, we do not accept the claim that the mixed-layer depth in a one-dimensional representation should be 700 m. Also, to correct the record, the main release of tritium was about 9 to 10 years (1962-1963) before the Atlantic GEOSECS (1972-1973) instead of Austin and Green’s 7 years. The use of a shorter time raised their $K$ estimate.

We would like to take this opportunity to correct a referencing error in our paper. The cruises GS 7307 and GS 7205 were USA cruises and the tritium data were obtained by the Tritium Laboratory of the University of Miami.